

## § 15. Possibility of Transport Barrier Formation by Bursting Toroidal Alfvén Eigenmodes

Toi, K., Yamamoto, S. (Dep. Energy Eng. Sci., Nagoya Univ.), Nakajima, N., Osakabe, M., Ohdachi, S., Sakakibara, S., Watanabe, K.Y., Narihara, K., Tanaka, K., Morita, S., Tokuzawa, T.

**Bursting toroidal Alfvén eigenmodes (TAEs) excited in NBI heated plasmas at low field  $B_t < 1\text{T}$  always induce enhanced radial transport and/or loss of energetic ions.** Figure 1 shows time evolution of the stored energy  $W_p$  measured in such plasma, where  $W_p$  is measured with a diaphragmatic loop. As seen from Fig.1,  $W_p$  has several small but sharp drops which synchronize with magnetic fluctuation burst related to  $n=2$  TAE. From the power balance for bulk plasma and beam ions generated by tangential NBI ( $W_p$  and  $W_{b//}$ ), the relative change of  $W_p$  is expressed as,

$$W_p/W_p(0) = \frac{\tau^*}{\tau_s} \frac{\tau^*(\tau_s - \tau^*)}{\tau_s(\tau_E - \tau^*)} \exp\left(-\frac{t}{\tau^*}\right) + \frac{\tau_E(\tau_s - \tau^*)}{\tau_s(\tau_E - \tau^*)} \exp\left(-\frac{t}{\tau_E}\right)$$

where  $\tau_s$ ,  $\tau_E$ , and  $\tau_c$  are respectively the slowing down time of beam ions, global energy confinement time and confinement time of energetic beam ion. The time  $\tau^*$  is expressed as

$$\tau^* = \frac{\tau_s \tau_c}{\tau_s + \tau_c}, \text{ and } \tau^* \approx \tau_c \text{ for } \tau_c \ll \tau_s. \text{ From the time}$$

evolution of  $W_p$  shown in Fig.1  $\tau_c$  is estimated to be  $\sim 3$  ms, for instance, at  $t \sim 1.35$  s. Therefore, the loss rate of energetic beam ions

$$(W_{b//}(0) - W_{b//})/W_{b//}(0) \approx 1 - \exp\left(-\frac{t}{\tau^*}\right) \text{ is estimated to}$$

be 30-50% for each TAE burst.

An interesting fact is found on the time evolution of  $W_p$  in Fig.1. That is,  $W_p$  rapidly recovers just after the TAE burst and even exceeds that just before the burst. Moreover, electron density and soft X-ray signals in the core region also increase slightly but clearly, synchronizing the  $W_p$ -increase, as shown in Fig.2[1]. This suggests the possibility of transport barrier formation induced by TAE burst. Transiently enhanced radial transport and/or loss of energetic ions as shown in Fig.1 may induce radial electric field shear and related poloidal shear flow. This process may be a trigger of transport barrier formation. A reason why the increase in  $W_p$  induced by TAE burst is fairly modest may be caused by the presence of rather large damping mechanisms of induced sheared flow. If the flow generation by frequent TAE burst prevails the damping, the rise in  $W_p$  by each

burst will pile up without the full decay [2]. A typical example of the case is shown in Fig.3.

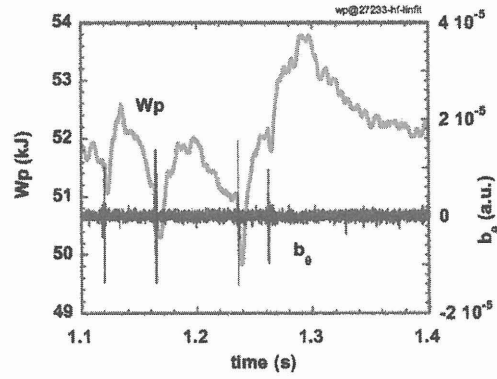


Fig.1 Time evolution of the stored energy  $W_p$  in an NBI heated discharge with TAE bursts.

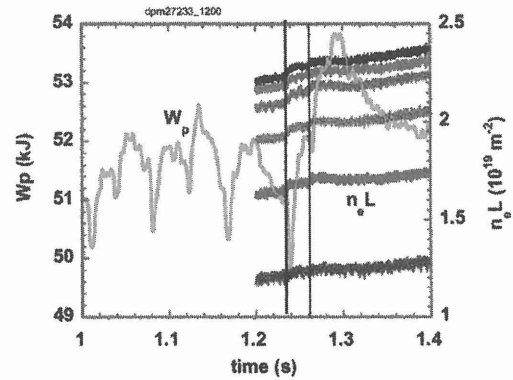


Fig.2 Time evolution of line-integrated electron density for various radial locations together with the bulk plasma energy.

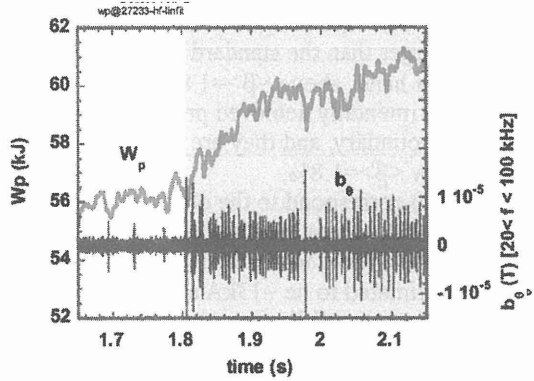


Fig.3 Time evolution of  $W_p$  in the phase that TAE burst is frequently excited.

## References

- [1] K. Toi et al., Proc. 19<sup>th</sup> IAEA Fusion Energy Conf., Lyon, 2002, paper No. EX/S-32.
- [2] GUENTER, S. et al., 28<sup>th</sup> EPS on Contr. Fusion and Plasma Phys. Funchal, 2001, No.P1-006.